Observation of $J/\psi \rightarrow p\bar{p}\alpha_0(980)$ at BESII

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I. INTRODUCTION

As one of the low-lying scalars, the state $a_0(980)$ has turned out to be mysterious in the quark model scenario. Its production near threshold allows tests of various hypotheses for its structure, including quark-antiquark [1], four quarks [2], $K \bar{K}$ molecule [3] and hybrid states [4]. The measurement of $J/\psi \rightarrow p\bar{p}a_0(980)$ is an additional observable constraining any phenomenological models trying to understand the nature of the $a_0(980)$.

A chiral unitary coupled channels approach of the chiral perturbation theory [5–7] is applied in investigation of the four-body decays $J/\psi \rightarrow N\bar{N}MM$ process [8] where the $N$ stands for a baryon and the $M$ for a meson. In this approach, the process $J/\psi \rightarrow p\bar{p}\pi^0\eta$ is investigated with the $a_0(980)$ meson generated through final state interaction (FSI). The amplitude of this process is calculable except for some coefficients which are not restricted, and its branching fraction varies within a wide range for different coefficients. Therefore, an experimental measurement of the process $J/\psi \rightarrow p\bar{p}a_0(980) \rightarrow p\bar{p}\pi^0\eta$ is needed for further progress in understanding the dynamics of the four-body decay processes taking the FSI of mesons into account.

In this paper, we present a measurement of $J/\psi \rightarrow p\bar{p}a_0(980)$ with $a_0(980)$ decaying to $\pi^0\eta$ based on $2.25 \times 10^8 J/\psi \rightarrow p\bar{p}a_0(980)$ events [9] collected with the BESIII detector at BEPCII.

II. THE EXPERIMENT AND DATA SETS

BESIII/BEPCII [10] is a major upgrade of BESII/BEPC [11]. BEPCII is a double-ring $e^+e^-$ collider running at 2.0–4.6 GeV center-of-mass energies; it provides a peak luminosity of $0.4 \times 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$ at the center-of-mass energy of 3.097 GeV.

The cylindrical BESIII detector has an effective geometrical acceptance of 93% of $4\pi$. It contains a small cell helium-based (40% He, 60% C$_2$H$_6$) main drift chamber (MDC) which has 43 cylindrical layers and provides an average single-hit resolution of 135 $\mu$m and momentum measurements of charged particles; a time-of-flight system (TOF) consisting of 5 cm thick plastic scintillators, with 176 detectors of length 2.4 m in two layers in the barrel and 96 fan-shaped detectors in the end caps; an electromagnetic calorimeter (EMC) consisting of 6240 CsI(Tl) crystals in a cylindrical structure and two end caps, which is used to measure the energies of photons and electrons; and a muon system consisting of resistive plate chambers. The momentum resolution of the charged particle is 0.5% at 1 GeV/c in a 1 tesla magnetic field. The energy loss $(dE/dx)$ measurement provided by the MDC has a resolution of 6%. The time resolution of the TOF is 80 ps in the barrel detector and 110 ps in the end cap detectors. The energy resolution of EMC is 2.5% (5.0%) in the barrel (end caps).

Monte Carlo (MC) simulated events are used to determine the detection efficiency, optimize selection criteria, and estimate possible backgrounds. The Geant4-based [12] simulation software Boost [13] includes the geometric and material description of the BESIII detectors, the detector response and digitization models, as well as the tracking of the detector running conditions and performance. The $J/\psi$ resonance is generated by KKMC [14] which is the event generator based on precise predictions of the electroweak standard model for the process $e^+e^- \rightarrow f\bar{f} + ny$, where $f = e, \mu, \tau, u, d, c, s, b$ and $n$ is an integer number $\geq 0$. The subsequent decays are generated with Evtden [15] with branching fractions being set to the world average values according to the Particle Data Group (PDG) [16] and the remaining unmeasured decays are generated by LUNDCHARM [17]. A sample of $2.25 \times 10^8$ simulated events, corresponding to the luminosity of data, is used to study background processes from $J/\psi$ decays (“inclusive backgrounds”). A signal MC sample with more than ten times the observed events in data for the process $J/\psi \rightarrow p\bar{p}a_0(980) \rightarrow p\bar{p}\pi^0\eta$ is generated, where the shape of the $a_0(980)$ is parametrized with the Flatté formula [18].

III. EVENT SELECTION

We select the process $J/\psi \rightarrow p\bar{p}\pi^0\eta$, with both $\pi^0$ and $\eta$ decaying to two photons, for this analysis. A good charged track is required to have good quality in the track fitting and be within the polar angle coverage of the MDC, i.e., $|\cos \theta| < 0.93$, and pass within 1 cm of the $e^+e^-$ interaction point in the transverse direction to the beam line and within 10 cm of the interaction point along the beam axis. Since the charged track in this process has relatively low transverse momentum, charged particle identification (PID) is only based on the $dE/dx$ information with the confidence level $\text{Prob}_{\text{PID}}(i)$ calculated for each particle hypothesis $i$ ($i = \pi/K/p$). A charged track with

Using $2.25 \times 10^8 J/\psi$ events collected with the BESIII detector at the BEPCII storage rings, we observe for the first time the process $J/\psi \rightarrow p\bar{p}a_0(980)$, $a_0(980) \rightarrow \pi^0\eta$ with a significance of $6.5\sigma$ (3.2$\sigma$ including systematic uncertainties). The product branching fraction of $J/\psi \rightarrow p\bar{p}a_0(980) \rightarrow p\bar{p}\pi^0\eta$ is measured to be $(6.8 \pm 1.2 \pm 1.3) \times 10^{-5}$, where the first error is statistical and the second is systematic. This measurement provides information on the $a_0$ production near threshold coupling to $p\bar{p}$ and improves the understanding of the dynamics of $J/\psi$ decays to four-body processes.
Prob_{pID}(p) > Prob_{pID}(K) and Prob_{pID}(p) > Prob_{pID}(\pi) is identified as a proton or an antiproton candidate. Photon candidates are required to have a minimum energy deposition of 25 MeV in the barrel (|cos \theta| < 0.8) of the EMC and 50 MeV in the end caps (0.86 < |cos \theta| < 0.92) of the EMC. EMC timing requirements (0 ≤ T ≤ 14 in units of 50 ns) are used to suppress electronic noise and to remove showers unrelated to the event. At the event selection level, candidate events are required to have at least two good charged tracks with one proton and one antiproton being identified, and at least four good photons.

We then perform a kinematic fit which imposes energy and momentum conservation at the production vertex to combinations of one proton and one antiproton candidate and four photons. For events with more than four photons, we consider all possible four-photon combinations, and the one giving the smallest χ^2 is the correct combination of photons. For events with four photons, the kinematic fit is selected for further analysis. To improve the signal-to-background ratio, events with χ^2_{4C} < 35 are accepted; this optimizes the figure of merit S/√S + B, where S and B are the numbers of MC simulated signal and inclusive background events respectively. The best photons pairing to π^0 and η in the four selected photons are selected by choosing the combination that gives the minimum χ^2-like variable

$$\chi^2_{\pi^0\eta} = \frac{(M_{\gamma\gamma\eta} - M_{\pi^0})^2}{\sigma_{\pi^0}^2} + \frac{(M_{\gamma\gamma\eta} - M_{\eta})^2}{\sigma_{\eta}^2},$$

where $M_{\gamma\gamma\eta}$ is the invariant mass of two photons after kinematic fit and $M_{\pi^0/\eta}$ is the π^0/η mass from PDG [16]. The mass resolutions for the π^0 and η, $\sigma_{\pi^0}$, and $\sigma_{\eta}$ are extracted by fitting the corresponding mass spectra in the signal MC sample; they are found to be 6.0 MeV/c^2 and 9.8 MeV/c^2 respectively. A MC study shows the rate of correct combination of photons is greater than 99% by using the $\chi^2_{\pi^0\eta}$ metric. To suppress p + π^0 + π^0 final states surviving in the 4C fit, we select two-photon pairs giving a minimum $\chi^2_{\pi^0\eta} = \frac{(M_{\gamma\gamma\eta} - M_{\pi^0})^2}{\sigma_{\pi^0}^2} + \frac{(M_{\gamma\gamma\eta} - M_{\eta})^2}{\sigma_{\eta}^2}$ and reject events with $\chi^2_{\pi^0\eta}$ less than 100. Figure 1 shows the mass spectra of selected $\gamma\gamma$ pairs for data and MC, where $\gamma_1\gamma_2$ indicates π^0 candidates and $\gamma_1\gamma_4$ indicates η candidates. The hatched histograms represent MC shapes from backgrounds and signal, where the background shapes are normalized based on their branching fractions and the signal shape is normalized to the rest area of the histogram of the data. We then require the mass of π^0 and η candidates to be within a 3σ window around their mean values.

IV. DATA ANALYSIS

The backgrounds contaminating the selected $J/\psi \rightarrow p\bar{p}\pi^0\eta$ candidates arise mainly from events with the same topology ($p\bar{p}\gamma\gamma\gamma\gamma$), events with an additional undetected photon ($p\bar{p}\gamma\gamma\gamma\gamma\gamma$), and events with a fake photon being reconstructed ($p\bar{p}\gamma\gamma\gamma\gamma$). The potential final states of background are categorized into four kinds: $p\bar{p}\pi^0\pi^0$, $p\bar{p}\pi^0\omega$, $p\bar{p}\eta\gamma$, and $p\bar{p}\pi^0\gamma\gamma$, where the π^0 can be produced from intermediate states Σ or Δ, and γ^0 can be produced from Ω. Since the branching fractions for the exclusive background processes $J/\psi \rightarrow \Sigma^+\Sigma^-\gamma/\Delta^+\Delta^-\gamma/\pi^0\omega(n\eta)$ have not yet been measured, we determine them from the same J/ψ data sample. The measurements are performed by requiring different numbers of photon candidates in one event and selecting the combination of π^0 with invariant mass closest to the mass of Σ or Δ, or selecting the combination of γ^0 closest to the mass of Ω. The measured branching fractions are shown in Table 1, where the uncertainty is statistical only. With the detection efficiency correction for the exclusive background satisfying the $p\bar{p}\pi^0\eta$ selection criteria, the contribution of the exclusive backgrounds is calculated to be 290 ± 19, which accounts for 4.3% of the surviving events found in data. The distributions of $M_{\pi^0/\eta}$ for data and backgrounds after normalization are presented in Fig. 2. A structure around 1.0 GeV [Fig. 2(a)] in data is clearly visible, but is not seen significantly in the
The likelihood method is used to fit the branching fraction of the processes with intermediate states of \( N^* \) in data. The maximum likelihood fit shows the scatter plot of the exclusive backgrounds in Fig. 2(b). The studies of the mass spectra of \( \psi \to \rho \eta \) described below. When applying the PWA without the component \( J/\psi \to \rho \eta \), no enhancement around 1.0 GeV is observed in the MC projection of \( \pi^0 \eta \) mass spectrum, which indicates that the enhancement seen in data is not from the processes with \( N^* \) intermediate states or their interferences.

An unbinned extended maximum likelihood fit is performed on the \( \pi^0 \eta \) mass spectrum. The probability density function is

\[
F(m) = f_{\text{sig}} \sigma(m) \otimes (e(m) \times \tilde{T}(m)) + (1 - f_{\text{sig}})B(m)
\]

Here, \( f_{\text{sig}} \) is the fraction of \( p \bar{p} a_0(980) \) signal events. The signal shape of \( a_0(980) \) is described as an efficiency-weighted Flatté formula \( (e(m) \times \tilde{T}(m)) \) convoluted with a resolution function \( \sigma(m) \). The non-\( a_0(980) \) background shape, expressed by \( B(m) \), is described by a third-order Chebychev polynomial function. The Flatté formula [18] is used to parametrize the \( a_0(980) \) amplitudes coupling to \( \pi^0 \eta \) and \( K \bar{K} \) by a two-channel resonance expressed as

\[
\tilde{T}(m) \propto \frac{1}{(m_{a_0}^2 - m^2)^2 + (\rho_{\pi^0\eta} g^2_{\omega\pi^0\eta} + \rho_{K\bar{K} a_0K\bar{K}} g^2_{a_0K\bar{K}})^2}
\]

### Table 1: Backgrounds of the final states with \( p \bar{p} \pi^0 \pi^0 \), \( p \bar{p} \pi^0 \pi^0 \), \( p \bar{p} \pi^0 \gamma \), and \( p \bar{p} \pi^0 \gamma \), where \( Br \) is the branching fraction of each channel, with statistical error only, \( \epsilon_{\text{sel}} \) is the efficiency of each channel determined with 50 k MC sample, and \( N_{\text{Norm}} \) is the number of background events normalized to the total \( J/\psi \) data.

<table>
<thead>
<tr>
<th>Channel ((J/\psi \to \to))</th>
<th>( Br )</th>
<th>( \epsilon_{\text{sel}} )</th>
<th>( N_{\text{Norm}} )</th>
</tr>
</thead>
<tbody>
<tr>
<td>( p \bar{p} \pi^0 \pi^0 )</td>
<td>((1.60 \pm 0.26) \times 10^{-3})</td>
<td>(1.68 \times 10^{-4})</td>
<td>(61 \pm 10)</td>
</tr>
<tr>
<td>( \Sigma^+ \Sigma^- \to p \pi^0 \bar{p} \pi^0 )</td>
<td>((2.77 \pm 0.03) \times 10^{-4})</td>
<td>(1.26 \times 10^{-4})</td>
<td>(8 \pm 0)</td>
</tr>
<tr>
<td>( \Delta^+ \Delta^- \to p \pi^0 \bar{p} \pi^0 )</td>
<td>((2.30 \pm 0.07) \times 10^{-4})</td>
<td>(1.76 \times 10^{-4})</td>
<td>(9 \pm 0)</td>
</tr>
<tr>
<td>( p \pi^0 \Delta^- \to p \pi^0 \bar{p} \pi^0 )</td>
<td>((2.04 \pm 0.06) \times 10^{-4})</td>
<td>(1.76 \times 10^{-4})</td>
<td>(8 \pm 0)</td>
</tr>
<tr>
<td>( \gamma \Sigma^+ \Sigma^- \to \gamma p \pi^0 \bar{p} \pi^0 )</td>
<td>((3.31 \pm 0.12) \times 10^{-5})</td>
<td>(2.98 \times 10^{-3})</td>
<td>(23 \pm 1)</td>
</tr>
<tr>
<td>( \gamma \Delta^+ \Delta^- \to \gamma p \pi^0 \bar{p} \pi^0 )</td>
<td>((5.40 \pm 0.50) \times 10^{-5})</td>
<td>(2.86 \times 10^{-3})</td>
<td>(35 \pm 3)</td>
</tr>
<tr>
<td>( \gamma p \pi^0 \Delta^- \to \gamma p \pi^0 \bar{p} \pi^0 )</td>
<td>((14.40 \pm 2.80) \times 10^{-5})</td>
<td>(2.44 \times 10^{-3})</td>
<td>(78 \pm 15)</td>
</tr>
<tr>
<td>( p \bar{p} \omega \to \gamma p \pi^0 \bar{p} \pi^0 )</td>
<td>((9.11 \pm 1.27) \times 10^{-5})</td>
<td>(1.59 \times 10^{-3})</td>
<td>(33 \pm 5)</td>
</tr>
<tr>
<td>( p \bar{p} \omega \to \gamma p \pi^0 \bar{p} \pi^0 )</td>
<td>((1.28 \pm 0.07) \times 10^{-5})</td>
<td>(1.14 \times 10^{-2})</td>
<td>(33 \pm 2)</td>
</tr>
<tr>
<td>( J/\psi \to \rho \eta )</td>
<td>((4.78 \pm 0.99) \times 10^{-7})</td>
<td>(1.80 \times 10^{-2})</td>
<td>(2 \pm 0)</td>
</tr>
<tr>
<td>Total</td>
<td></td>
<td></td>
<td>(290 \pm 19)</td>
</tr>
</tbody>
</table>

FIG. 2 (color online). (a) The mass spectrum of \( \pi^0 \eta \) for data and exclusive backgrounds. The dots with error bars represent data and the others are exclusive backgrounds after normalization. (b) The mass spectra of \( \pi^0 \eta \) for exclusive backgrounds.
where \( \rho_{\eta^0}\) and \( \rho_{KK}\) are the decay momenta of the \( \pi^0 \) or \( K \) in the \( \pi^0\eta \) or \( KK \) rest frame, respectively. The two coupling constants \( g_{a_0\pi^0}\) and \( g_{a_0KK}\) stand for \( a_0(980) \) resonance coupling to \( \pi^0\eta \) and \( KK \), respectively. The experiment results from Refs. [20–22] are consistent with each other and the weighted average of them is calculated as \( g_{a_0\pi^0} = 2.83 \pm 0.05 \) and \( g_{a_0KK} = 2.11 \pm 0.06 \). In the fit, the two coupling constants \( g_{a_0\pi^0}\) and \( g_{a_0KK}\) are fixed to 2.83 and 2.11, respectively.

The mass-dependent efficiency \( \epsilon(m) \) is studied by using a large phase space MC \( J/\psi \to p\bar{p}\pi^0\eta \) sample, where the efficiency curve derived from the four-body phase space MC is compatible with that from signal MC of \( p\bar{p}a_0(980) \). The detector resolution \( \sigma(m) \) of \( M_{\pi^0\eta} \) is extracted by using a large sample of simulated signal events \( J/\psi \to p\bar{p}a_0(980) \), \( a_0(980) \to \pi^0\eta \), with the width of the \( a_0(980) \) set to zero.

In the fit, the signal fraction \( f_{\text{sig}} \), the \( a_0(980) \) mass, and the parameters of the background polynomial are allowed to vary. The fit result of \( M_{\pi^0\eta} \) is shown in Fig. 4. The yield of \( a_0(980) \) events is \( 849 \pm 144 \), with a statistical significance of \( 6.5\sigma \) which is calculated from the log-likelihood difference between fits with and without the \( a_0(980) \) signal component. The fit mass is \( 1.012 \pm 0.007 \text{ GeV}/c^2 \), which is slightly higher than the PDG value [16]. The robustness of this result has been validated with a toy MC study. Different signal MC samples of \( J/\psi \to p\bar{p}a_0(980) \), \( a_0(980) \to \pi^0\eta \) are generated with different mass and width of the \( a_0(980) \). Background events are randomly sampled according to the background shapes. In all cases, the fit value of the \( a_0(980) \) mass is found to be consistent with the input value within statistical uncertainties. The product branching fraction \( \text{Br}(J/\psi \to p\bar{p}a_0(980) \to p\bar{p}\pi^0\eta) \) is calculated to be \( (6.8 \pm 1.2) \times 10^{-5} \), where the uncertainty is statistical only.

**V. ESTIMATION OF SYSTEMATIC UNCERTAINTIES**

The systematic uncertainties on the measurement of \( \text{Br}(J/\psi \to p\bar{p}a_0(980) \to p\bar{p}\pi^0\eta) \) are summarized in Table II. Systematic uncertainties due to tracking and PID efficiency, photon detection efficiency, the kinematic fit, and the \( \pi^0\eta \) veto arise due to imperfect modeling of the data by the simulation. The systematic uncertainty associated with the tracking efficiency as a function of transverse momentum and the uncertainty due to the PID efficiency of proton/antiproton have been studied by a control sample of \( J/\psi \to p\bar{p}\pi^+\pi^- \) decays using a technique similar to that discussed in Ref. [23]. In this paper, due to the low transverse momentum of proton and antiproton, the uncertainty of tracking efficiency is determined by the weighted uncertainty \( \Sigma_i \epsilon_i r_i \), where \( \epsilon_i \) represents the data/MC difference in each transverse momentum bin [23] and \( r_i \) represents the proportion of each transverse momentum bin in data. The systematic uncertainty due to the tracking efficiency is estimated to be 4.0% per proton and 5.0% per antiproton, respectively. The large uncertainty of tracking

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**FIG. 3** (color online). (a) The scatter plot of \( M_{p\pi^+}^2 \) versus \( M_{\eta^0}^2 \) from data. (b) The scatter plot of \( M_{p\pi^+}^2 \) versus \( M_{\eta^0}^2 \) from MC projection of all intermediate states superimposed.

**FIG. 4** (color online). The results of fitting the mass spectrum for \( \pi^0\eta \). Dots with error bars are data and the solid line is the fitted spectrum. The dash-dotted line shows the non-\( a_0 \) background described by a third-order Chebychev polynomial. The dashed line shows the signal described by an efficiency-weighted Flatté formula convoluted with a resolution function.
errors are statistical only. Conservatively, the systematic uncertainty associated with the kinematic fit is determined by the difference of efficiencies for MC and data, where the first uncertainty is statistical and the second is systematic. Without considering the interference between the signal channel and the same final states with intermediate $N^+$ states, the branching fraction is measured to be

$$\text{Br}(J/\psi \to p\bar{p}a_0(980)) = (6.8 \pm 1.2 \pm 1.3) \times 10^{-5},$$

where the first uncertainty is statistical and the second is systematic.

Our measurement provides a quantitative comparison with the chiral unitary approach [8]. This approximation uses several coefficients in the parametrization of meson-meson amplitudes. One of them, namely $r_4$ in [8], is constrained by fitting the $\pi^+\pi^-$ invariant mass distribution in the decay $J/\psi \to p\bar{p}\pi^0\pi^0$; the fits suggests two equally possible values, $r_4 = 0.2$ and $r_4 = -0.27$. The theory also predicts that the branching fractions of $J/\psi \to p\bar{p}a_0(980)$ and $J/\psi \to p\bar{p}\pi^0\pi^0$ are comparable for $r_4 = -0.27$, while the branching fraction of the former is one or two orders of magnitude lower than that of the latter for $r_4 = 0.2$. Taking the branching fraction of $J/\psi \to p\bar{p}\pi^0\pi^0$ from PDG [16], the ratio of $\text{Br}(J/\psi \to p\bar{p}a_0(980))$ to $\text{Br}(J/\psi \to p\bar{p}\pi^0\pi^0)$ is found to be about $10^{-2}$, which shows preference to $r_4 = 0.2$.

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